

**Erratum: High-field magnetoconductivity of electrons on hydrogen  
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The expression for  $\sigma_{xx}$  appearing in Eq. (2) is in error. The correct expression includes a contribution from both spin species, and can be written as

$$\sigma_{xx} = \frac{\alpha e^2}{2\pi\hbar} \int_{-\infty}^{\infty} dE \left( \frac{-\partial f}{\partial E} \right) \sum_{N=0}^{\infty} \left( N + \frac{1}{2} \right) \left[ \exp \left( \frac{-4(E - E_N - \Delta S)^2}{\Gamma^2} \right) + \exp \left( \frac{-4(E - E_N + \Delta S)^2}{\Gamma^2} \right) \right],$$

where  $\Delta S = g\mu_B B/2$  and  $\alpha$  is an adjustable parameter. Theoretically  $\alpha$  should be unity but we obtained reasonably good fits with  $\alpha \approx 2$ . The origin of this discrepancy, which was also observed by Van De Sanden *et al.* (Ref. 8), is not understood.

The density of states used in the above expression for  $\sigma_{xx}$  is

$$D(E) = \frac{1}{2\pi l_0^2} \left( \frac{2}{\pi} \right)^{1/2} \frac{1}{\Gamma} \sum_{N=0}^{\infty} \left[ \exp \left( \frac{-2(E - E_N - \Delta S)^2}{\Gamma^2} \right) + \exp \left( \frac{-2(E - E_N + \Delta S)^2}{\Gamma^2} \right) \right],$$

which is equivalent to that quoted in Eq. (4). However, one can easily show using the above expression that, after normalizing,  $\sigma_{xx}$  is independent of  $\Delta S$  and the Zeeman splitting plays no role in the transport.

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